### Phonon dispersion relations for Silicon

### <u>Note:</u>

- LA (compression wave) alsways higher frequency than TA (shear wave)
- LO + 2 x TO degenerate at Γ<sup>+</sup>
- 2xTO stay degenerate along the symmetry lines
   Δ and Λ.

For q along high-symmetry directions, modes can be classified as



Accoustic modes







C not only lighter than Si, but in addition, C-C-bonds much stronger than Si-Si bonds!

## Comparison: Silicon and GaAs



<u>Si:</u>

#### <u>GaAs:</u>

 $M_{Si}$ =28 u

LO/TO degenerate at  $\Gamma^{\scriptscriptstyle +}$ 

M= 68 u = 1,5<sup>2</sup> M<sub>Si</sub>

LO/TO split at  $\Gamma^+$  due to inequivalent atoms in the base!



- Larger frequency (=energy) gap between optical and acoustic branches
- > Otical branches get flatter!

Already reproduced by the ,linear chain model' (see, e.g. Ibach)

# Phonon Density of States



\*Debye approximation for the phonons :

Frequency

Three phonon branches with linear dispersion relations  $\omega_n(\vec{q}) = v_n \cdot |\vec{q}|$  corresponding to the low-frequency ranges of the TA and LA branches and a common cut-off frequency such that the correct total number of phonon modes (including the optical ones) is reproduced.

$$N_{n}(\omega) = \frac{1}{(2\pi)^{3}} \cdot \frac{4}{3} \pi \left(\frac{\omega}{v_{n}}\right)^{3} \longrightarrow Z_{n}(\omega) = \frac{dN_{n}}{d\omega} = \frac{4\pi}{(2\pi)^{3} v_{n}^{3}} \cdot \omega^{2}$$
and thus for n = TA, TA, TO
$$Z(\omega) = \frac{4\pi}{(2\pi)^{3}} \cdot \left(\frac{2}{v_{TA}^{3}} + \frac{1}{v_{LA}^{3}}\right) \cdot \omega^{2}$$